

The Crystal Structure of Aflatoxin B₂

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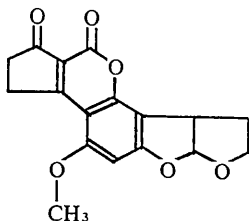
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Aflatoxin B₂ crystallizes as orthorhombic bipyramids in the space group $P2_12_12_1$ with four molecules in the unit cell with dimensions $a=6.927$, $b=10.146$, $c=20.171$ Å. The structure, solved by Patterson methods, consists of strings of coplanar molecules in the direction of the b axis. The molecules within a string are interconnected by an interaction between the methyl group and a carbonyl oxygen atom, which may be called a hydrogen bond; the $\text{CH}_3 \cdots \text{O}$ distance is 2.930 Å. The strings are stacked at an interplanar distance of 3.463 Å in the direction of the a axis. The libration tensor of the molecule is practically diagonal in the reference system defined by the principal axes of the inertia tensor.

Introduction

Aflatoxin B₂ is one of the carcinogenic metabolites of the mould *Aspergillus flavus* Link ex Fries (Hartley, Nesbitt & O'Kelly, 1963), which, in 1960 caused the death of 100,000 young turkeys as a result of liver damage (turkey X-disease). An excellent review of aflatoxins, including for instance carcinogenic action, metabolism and biochemical effects on DNA, RNA and proteins, has been presented by Schoental (1967).

According to van Dorp, van der Zijden, Beerthuis, Sparreboom, Ord, de Jong & Keuning (1963), aflatoxin B₂ has the following structural formula:



This formula has been confirmed by van Soest & Peerdeman (1964).

The present paper describes in detail the crystal and molecular structure of aflatoxin B₂.

Experimental

Aflatoxin B₂ (kindly supplied by Professor D. A. van Dorp of the Unilever Research Laboratory, Vlaardingen) was recrystallized from a chloroform-ethanol mixture. The crystals obtained were pale-yellow, well-formed orthorhombic bipyramids.

The unit-cell dimensions were first determined from Weissenberg and rotation photographs and refined,

later on, by a least-squares fit of 60 θ -values ($\theta > 68^\circ$) measured with a General Electric diffractometer (Furnas, 1957) using Cu radiation, $\lambda(K\alpha_1) = 1.54051$ Å.

The space group was uniquely determined from the systematically absent reflexions.

Crystal data

Aflatoxin B₂, $\text{C}_{17}\text{H}_{14}\text{O}_6$ ($M = 314.3$). Orthorhombic $a = 6.927 \pm 0.002$, $b = 10.146 \pm 0.002$, $c = 20.171 \pm 0.002$ Å. $U = 1417.7$ Å³, $D_m = 1.47$ g.cm⁻³, $Z = 4$, $D_x = 1.472$ g.cm⁻³. Space group $P2_12_12_1$, $\mu = 11.0$ cm⁻¹ for Cu $K\alpha$ radiation.

Two sets of intensity values were measured with a General Electric diffractometer, using nickel-filtered Cu radiation and a scintillation counter with pulse height discrimination. By measuring the apparent absorption factor for a nickel filter as a function of the intensity of the primary beam, it was found that the counting apparatus was linear up to only 1000 counts.sec⁻¹. Seven attenuating filters of increasing thickness were used to keep the counting rate below this value. These intensities were measured by the θ - 2θ scanning technique at a speed of 2°.min⁻¹. The scan range was adjusted to account for the α_1 - α_2 splitting. Background was counted for 15 sec at each end of the scan. 1106 of the 1895 theoretically possible independent reflexions were measured, using a crystal with a largest dimension of 0.2 mm. For a second set of data (1725 reflexions measured), a different crystal was used with a maximum dimension of 0.6 mm. The intensities were corrected for Lorentz and polarization factors as well as for absorption – using a program written by Duisenberg (1966) by which absorption corrections for a crystal of arbitrary shape can be calculated – and placed on an approximately absolute scale by Wilson's method.

Table 1. Observed and calculated structure factors (x 10) and phase angles. Asterisks indicate unobserved reflexions. E's denote extinction.

Table with multiple columns of numerical data representing structure factors and phase angles for various hkl reflections. The table is organized into several vertical sections, each with its own set of column headers. Asterisks and 'E' symbols are used to denote specific data points.

Determination and refinement of the structure

A three-dimensional vector map, sharpened by the method of Jacobson, Wunderlich & Lipscomb (1961) was calculated. In this vector map, layers of high density are visible with maxima at the sections $u=0.000$, 0.132 , 0.368 and 0.500 (Fig. 1). First, this means that the aflatoxin B₂ molecules have a plane structure and are parallel to the (100) plane and, second, that the distances from the molecule to two of the three screw axes are $0.066 x/a$ and $0.184 x/a$. It is clear from the characteristic distribution of the peaks in these sections that the molecule contains a number of hexagonal rings. The oval peaks near the origin in the section with $u=0.000$ indicate that the rings are not altogether oriented parallel to the b axis. From these peaks, a rotation of about $\pm 4^\circ$ from the axis can be deduced.

During the work on this structure, the structural formula of aflatoxin B₁ became known from a paper by Asao, Büchi, Abdel-Kader, Chang, Wick & Wogan (1963), from which formula we could derive a plane model containing 19 atoms. (At that stage we did not know that the carbon atom of the methoxyl group was also in the same plane.)

Three vector maps, containing the endpoints of vectors between trial molecules interrelated by screw axes, were drawn on transparent paper and used to analyse the three sections containing the intermolecular vectors. This resulted in the correct position of the model with respect to the screw axes. The four missing atoms were subsequently found by Fourier methods. Initial refinement was done with 412 structure factors ($|F| > 10$) of the first set of data and described by van Soest & Peerdeman (1964). A further refinement was carried out on Telefunken TR4 and Electrologica X-8 computers with a block-diagonal least-squares program written by Palm & Peterse (1964). The second set of data was used after omitting 56 reflexions supposed to be affected by extinction (found by comparison of the two sets of data) and replacement of the non-observed reflexions by half the minimum value observed in corresponding ranges of $\sin \theta/\lambda$.

The refinement was started at an overall temperature factor of 2.7 \AA^2 . During five cycles, $R = \sum(|F_o| - |F_c|) / \sum|F_o|$ fell to 0.18. Four further cycles with individual isotropic temperature factors reduced R to 0.142 and three more with anisotropic temperature factors to 0.081. At this point, a difference map was calculated revealing the approximate positions of the 14 hydrogen atoms. The hydrogen atoms were given isotropic temperature factors equal to those of the carbon atoms to which they are bound. During the next four cycles, R dropped to 0.051; the positions of all the atoms, the anisotropic temperature factors of the heavy atoms and the scaling factor were refined. Using these parameters, the 56 structure factors of the large crystal supposed to be affected by extinction were calculated, giving an R factor of 0.20 for these structure factors, whereas for the corresponding structure factors of the small crystal

$R=0.07$. For 8 of these 56 reflexions, F_c is still greater than F_o and after omitting them, R decreased from 0.07 to 0.03.

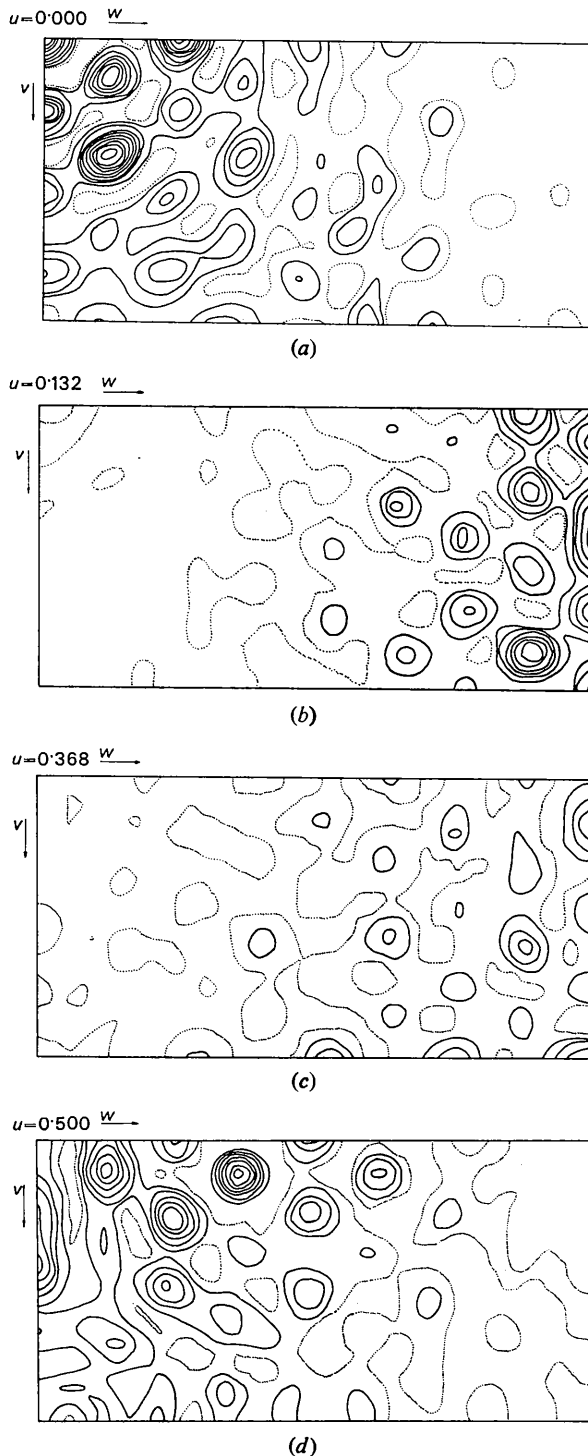


Fig. 1. Sections through the sharpened Patterson function $P(u, v, w)$ for the values u (a) 0.000; (b) 0.132; (c) 0.368 and (d) 0.500. Contours are at arbitrary intervals: solid, positive; dotted, zero.

The 48 remaining structure factors of the small crystal were included in the last three cycles, using weights inversely proportional to mean values of $\sum(|F_o| - |F_c|)^2$ in corresponding ranges of $|F_o|$ as calculated during the refinement. After the last cycle, the shifts in the positional and thermal parameters for non-hydrogen atoms were all less than $\sigma/4$ and $\sigma/3$ respectively. Those for the hydrogen atoms were less than $\sigma/2$. The final R for 1,717 reflexions was 0.046. A list of the final structure factors [here and in the preceding cases calculated by using the atomic scattering factors taken from *International Tables for X-ray Crystallography* (1962)] is given in Table 1 and the structural parameters are given in Tables 2 to 4. The arbitrary numbering of the atoms of the molecule is given in Fig. 2.

Table 2. Final positional parameters and *e.s.d.*'s ($\times 10^5$) for the carbon and oxygen atoms

	<i>x/a</i>	<i>y/b</i>	<i>z/c</i>
C(1)	-6577 (41)	72061 (21)	69511 (11)
C(2)	-6796 (39)	61294 (23)	64294 (11)
C(3)	-6477 (38)	48191 (21)	68027 (11)
C(4)	-6624 (37)	65511 (20)	75993 (11)
C(5)	-6542 (34)	52040 (20)	75277 (11)
C(6)	-7064 (41)	71954 (21)	82353 (11)
C(7)	-6500 (37)	43694 (20)	80971 (11)
C(8)	-6609 (38)	29701 (21)	80780 (11)
C(9)	-6717 (36)	49929 (20)	87206 (11)
C(10)	-6794 (43)	22229 (22)	86568 (12)
C(11)	-6872 (42)	42764 (24)	93011 (12)
C(12)	-7071 (41)	29161 (23)	92469 (12)
C(13)	-6638 (53)	46792 (28)	180 (12)
C(14)	-6801 (63)	33168 (31)	3665 (13)
C(15)	-6179 (56)	10448 (23)	73928 (15)
C(16)	+12162 (66)	52855 (33)	2592 (14)
C(17)	+24074 (68)	40972 (42)	4541 (17)
O(1)	-6425 (44)	83744 (17)	68374 (9)

Table 2 (cont.)

	<i>x/a</i>	<i>y/b</i>	<i>z/c</i>
O(2)	-7760 (37)	83462 (16)	83581 (9)
O(3)	-6908 (29)	63472 (15)	87826 (7)
O(4)	-7516 (38)	23089 (18)	98465 (9)
O(5)	-6167 (34)	24503 (15)	74679 (8)
O(6)	+10166 (47)	31964 (24)	7240 (10)

Table 3. Final positional parameters and *e.s.d.*'s ($\times 10^4$) and isotropic temperature factors (\AA^2) for the hydrogen atoms

	<i>x/a</i>	<i>y/b</i>	<i>z/c</i>	<i>B</i>
H(1)	-1936 (51)	6254 (34)	6178 (17)	2.37
H(2)	+505 (52)	6190 (33)	6168 (16)	2.37
H(3)	-1739 (50)	4254 (33)	6703 (16)	2.11
H(4)	+467 (53)	4299 (31)	6702 (15)	2.11
H(5)	-572 (61)	943 (37)	6907 (18)	3.83
H(6)	-1567 (58)	701 (36)	7615 (18)	3.83
H(7)	+351 (58)	736 (35)	7611 (18)	3.83
H(8)	-773 (56)	1297 (32)	8637 (16)	2.69
H(9)	-1701 (55)	5078 (36)	100 (19)	3.56
H(10)	-1690 (63)	3071 (41)	729 (18)	4.32
H(11)	+1539 (60)	5751 (40)	9936 (19)	4.78
H(12)	+739 (67)	5825 (37)	613 (19)	4.78
H(13)	+3035 (64)	3707 (42)	76 (21)	5.67
H(14)	+3064 (70)	4260 (44)	827 (21)	5.67

Discussion

The bond lengths and the bond angles (not corrected for thermal motion) together with their standard deviations are listed in Tables 5 and 6. The bond lengths are also shown in Fig. 2. Aflatoxin B₂ is largely a plane molecule: all heavy atoms but C(16), C(17), and O(6) lie within 0.05 Å from a plane given by the least-squares

Table 4. Final thermal parameters and *e.s.d.*'s ($\times 10^5$) for the carbon and oxygen atoms

The anisotropic temperature factor is defined as:

$$\exp [-(\beta_{11}h^2 + \beta_{22}k^2 + \beta_{33}l^2 + 2\beta_{12}hk + 2\beta_{13}hl + 2\beta_{23}kl)] .$$

	β_{11}	β_{22}	β_{33}	β_{12}	β_{13}	β_{23}
C(1)	1463 (49)	489 (19)	159 (5)	+49 (32)	-16 (16)	+24 (8)
C(2)	1271 (45)	614 (20)	126 (5)	+6 (34)	+13 (14)	+16 (8)
C(3)	1123 (42)	504 (18)	123 (5)	-16 (39)	-24 (14)	-15 (8)
C(4)	1188 (43)	465 (18)	139 (5)	+6 (29)	-6 (14)	-9 (8)
C(5)	865 (38)	437 (18)	132 (5)	+38 (27)	+12 (13)	-18 (7)
C(6)	1431 (47)	487 (19)	148 (5)	-29 (31)	-58 (16)	-6 (8)
C(7)	1043 (42)	455 (18)	135 (5)	+8 (29)	-14 (15)	+22 (7)
C(8)	1119 (42)	484 (19)	165 (5)	-7 (39)	-40 (15)	+1 (8)
C(9)	1000 (41)	483 (19)	135 (5)	+41 (29)	-5 (14)	+12 (8)
C(10)	1527 (50)	486 (20)	188 (6)	-124 (33)	-19 (18)	+50 (9)
C(11)	1348 (48)	678 (22)	140 (5)	-7 (34)	+56 (17)	+9 (9)
C(12)	1349 (47)	660 (22)	154 (5)	-99 (34)	+2 (16)	+102 (9)
C(13)	2548 (74)	865 (26)	131 (5)	+188 (49)	+134 (21)	+19 (10)
C(14)	3333 (93)	1007 (31)	143 (6)	+108 (59)	+156 (22)	+87 (11)
C(15)	2990 (80)	404 (20)	246 (7)	+13 (42)	-6 (25)	-44 (10)
C(16)	4007 (120)	1028 (33)	144 (6)	-192 (57)	-168 (24)	-12 (12)
C(17)	3626 (122)	1545 (50)	221 (9)	+16 (77)	-361 (28)	+28 (17)
O(1)	4047 (72)	497 (16)	195 (5)	-14 (36)	-24 (18)	+70 (7)
O(2)	2927 (55)	435 (14)	190 (4)	+67 (30)	-43 (15)	-50 (7)
O(3)	1794 (37)	455 (14)	127 (4)	+96 (24)	+16 (11)	-31 (5)
O(4)	2881 (56)	784 (18)	167 (4)	-179 (34)	+67 (15)	+133 (7)
O(5)	2760 (52)	393 (13)	161 (4)	+14 (28)	-52 (14)	-24 (6)
O(6)	4155 (83)	1205 (26)	178 (5)	+47 (45)	-181 (18)	+118 (10)

equation

$$-0.99996X - 0.00552Y - 0.00764Z - 0.3139 = 0$$

The tetrahydrofuran ring protrudes from this plane and the equation of the plane formed by the four atoms C(13), C(14), C(16) and O(6) is

$$-0.46563X + 0.38972Y + 0.79455Z - 18.1320 = 0$$

Table 5. Bond lengths and *e.s.d.*'s (Å)

C-C bonds		- C-O bonds	
<i>sp</i> ³ - <i>sp</i> ³			
C(2)-C(3)	1.528 (3)	C(14)-O(6)	1.385 (5)
C(13)-C(14)	1.553 (4)	C(15)-O(5)	1.437 (3)
C(13)-C(16)	1.523 (5)	C(14)-O(4)	1.465 (3)
C(16)-C(17)	1.513 (6)	C(17)-O(6)	1.434 (5)
<i>sp</i> ² - <i>sp</i> ³			
C(1)-C(2)	1.521 (3)	= C-O bonds	
C(3)-C(5)	1.516 (3)	C(9)-O(3)	1.381 (3)
C(11)-C(13)	1.504 (3)	C(8)-O(5)	1.339 (3)
<i>sp</i> ² - <i>sp</i> ²			
C(1)-C(4)	1.467 (3)	C(12)-O(4)	1.359 (3)
C(4)-C(6)	1.437 (3)	C(6)-O(3)	1.401 (3)
C(5)-C(7)	1.426 (3)	C-H bonds	
C=C bonds			
C(4)-C(5)	1.377 (3)	C(2)-H(1)	0.98 (4)
C...C bonds			
C(7)-C(8)	1.424 (3)	C(2)-H(2)	1.01 (4)
C(7)-C(9)	1.407 (3)	C(3)-H(3)	0.96 (4)
C(8)-C(10)	1.390 (3)	C(3)-H(4)	0.97 (4)
C(9)-C(11)	1.377 (3)	C(10)-H(8)	0.94 (4)
C(10)-C(12)	1.383 (3)	C(13)-H(9)	0.84 (4)
C(11)-C(12)	1.382 (3)	C(14)-H(10)	1.04 (4)
C=O bonds			
C(1)-O(1)	1.206 (3)	C(15)-H(5)	0.99 (4)
C(6)-O(2)	1.197 (3)	C(15)-H(6)	0.87 (4)
		C(15)-H(7)	0.86 (4)
		C(16)-H(11)	0.84 (4)
		C(16)-H(12)	0.96 (4)
		C(17)-H(13)	0.96 (4)
		C(17)-H(14)	0.89 (4)

Table 6. Bond angles and *e.s.d.*'s for the non-hydrogen atoms

C(2)-C(1)-O(1)	125.08 (0.16)
C(4)-C(1)-O(1)	127.96 (0.20)
C(2)-C(1)-C(4)	106.94 (0.24)
C(1)-C(2)-C(3)	106.52 (0.15)
C(2)-C(3)-C(5)	104.60 (0.22)
C(1)-C(4)-C(5)	110.94 (0.17)
C(1)-C(4)-C(6)	126.17 (0.25)
C(5)-C(4)-C(6)	122.87 (0.18)
C(3)-C(5)-C(4)	110.95 (0.18)
C(3)-C(5)-C(7)	128.74 (0.23)
C(4)-C(6)-O(2)	128.90 (0.20)
C(4)-C(6)-O(3)	115.30 (0.24)
O(2)-C(6)-O(3)	115.78 (0.15)
C(5)-C(7)-C(8)	124.79 (0.16)
C(5)-C(7)-C(9)	117.15 (0.24)
C(8)-C(7)-C(9)	118.04 (0.18)
C(7)-C(8)-O(5)	114.63 (0.18)
C(7)-C(8)-C(10)	121.38 (0.18)
C(10)-C(8)-O(5)	123.96 (0.25)
C(8)-O(5)-C(15)	119.14 (0.17)
C(6)-O(3)-C(9)	122.60 (0.14)
C(7)-C(9)-O(3)	121.71 (0.18)
C(7)-C(9)-C(11)	121.65 (0.26)
O(3)-C(9)-C(11)	116.60 (0.18)

Table 6 (cont.)

C(8)-C(10)-C(12)	116.50 (0.28)
C(9)-C(11)-C(12)	117.28 (0.19)
C(9)-C(11)-C(13)	132.24 (0.31)
C(12)-C(11)-C(13)	110.46 (0.23)
C(10)-C(12)-C(11)	125.11 (0.21)
C(10)-C(12)-O(4)	122.41 (0.31)
C(11)-C(12)-O(4)	112.46 (0.20)
C(11)-C(13)-C(14)	101.00 (0.29)
C(11)-C(13)-C(16)	115.38 (0.34)
C(14)-C(13)-C(16)	102.87 (0.38)
C(13)-C(14)-O(4)	107.32 (0.18)
C(13)-C(14)-O(6)	107.74 (0.45)
O(4)-C(14)-O(6)	110.00 (0.45)
C(13)-C(16)-C(17)	103.12 (0.48)
C(16)-C(17)-O(6)	103.86 (0.56)
C(14)-O(6)-C(17)	108.58 (0.36)
C(12)-O(4)-C(14)	107.26 (0.28)

Table 7. Distances from the heavy atoms to the least-squares planes of the aflatoxin B₂ molecule (Å)

$$-0.99996X - 0.00552Y - 0.00764Z - 0.3139 = 0$$

C(1)	-0.007
C(2)	+0.023
C(3)	+0.003
C(4)	-0.010
C(5)	-0.006
C(6)	+0.007
C(7)	-0.013
C(8)	+0.003
C(9)	-0.011
C(10)	+0.012
C(11)	-0.005
C(12)	+0.018
C(13)	-0.035
C(14)	-0.021
C(15)	-0.004
O(1)	-0.023
O(2)	+0.047
O(3)	-0.007
O(4)	+0.043
O(5)	-0.015

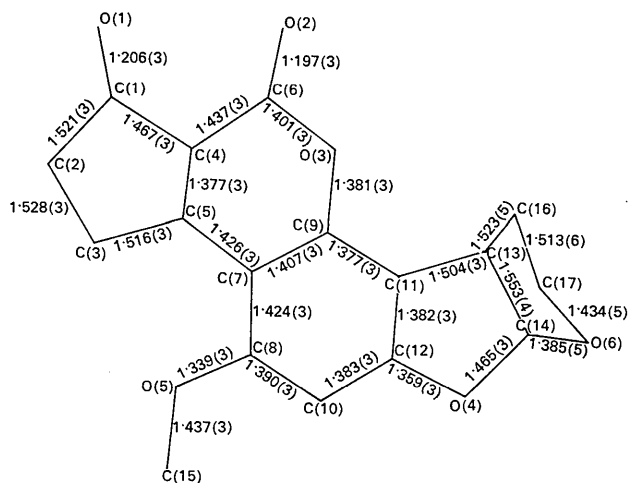


Fig. 2. Arbitrary numbering of the atoms, the bond lengths and their standard deviations of the aflatoxin B₂ molecule.

Table 7 (cont.)

$$-0.46563X + 0.38972Y + 0.79455Z - 18.1220 = 0$$

C(13)	-0.012
C(14)	+0.013
C(16)	+0.008
O(6)	-0.009
[C(17)]	-0.534

The angle between these two planes is 117.21° . The distances from the atoms to their respective planes are listed in Table 7. The bond between the carbon atoms with the highest degree of substitution C(13) and C(14) is significantly longer than the other three sp^3-sp^3 bonds, as has been observed *e.g.* for iso-eremolactone (Oh & Maslen, 1968). The value of 1.467 \AA for the sp^2-sp^2 bond in the cyclopentenone ring agrees well with 1.462 \AA in cyclo-octatetraene (Bastiansen, Hedberg & Hedberg, 1957), whereas the two sp^2-sp^2 bonds in the lactone ring are significantly shorter. Although the variations in the six bond lengths of the benzene nucleus are considerable, the mean value of 1.394 \AA is in agreement with the value of 1.394 \AA given by Sutton (1965). Among the $-C-O$ bonds, the short C(14)-O(6) bond of 1.385 \AA deserves attention for it causes an asymmetry of the C(14)-O(4) and C(14)-O(6) bonds.

The $=C-O$ bonds present in a ring system are all longer than 1.339 \AA , which is the length of the $=C(8)-O(5)$ bond in the methoxyl group. Oh & Maslen (1968) have ascribed the extension of the two $=C-O$ bonds (1.379 and 1.400 \AA) in the γ -lactone ring of iso-eremolactone to the small interior angles that would induce a greater p character in the bonds. However, in the δ -lactone ring of aflatoxin B_2 , the angles are normal and the same extension of the bond lengths occurs. The difference in the carbonyl bond angles O(2)-C(6)-C(4) and O(2)-C(6)-O(3) seems to be normal for γ -lactone rings [*cf.* iso-eremolactone and α -methyltetronic acid (MacDonald & Alleyne, 1963)] and for δ -lactone rings [*cf.* 4-hydroxycoumarin (Gaultier & Hauw, 1966) and 3-bromo-4-hydroxycoumarin monohydrate (Gaultier & Hauw, 1965)]. The smaller angle of the two is always at the side of the oxygen atom in the ring. The methoxyl group is coplanar with respect to the benzene nucleus of the molecule, as has been found *e.g.* for 2-*p*-methoxyphenyl-3,4-dibenzyl-1,3,4-thiadiazolidine-5-thione (Karle & Karle, 1965) and in myxin (Hanson, 1968*b*). However, in dibromo-erioistic acid (Paton, Maslen & Watson, 1967), methyl is protruding from the plane. In aflatoxin B_2 as well as in myxin and dibromoerioistic acid, the methyl hydrogen atoms have a staggered con-

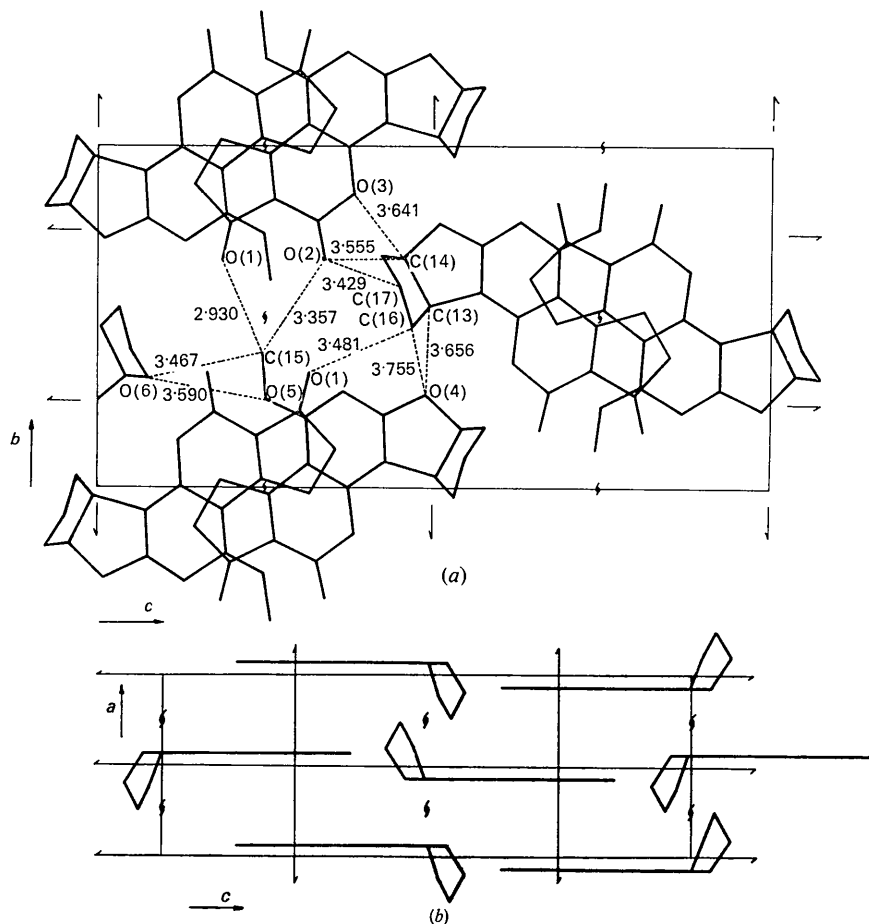


Fig. 3. Projection of the structure of aflatoxin B_2 (a) on the (100) plane and (b) on the (010) plane.

formation with respect to the C(benzene nucleus)–O bond. For this conformation, there is also little steric hindrance from the hydrogen atom of the benzene nucleus. The carbon–hydrogen bond lengths vary from 0.84 to 1.04 Å with a mean value of 0.94 Å and a standard deviation of 0.06 Å.

The crystal structure of aflatoxin B₂ is shown in Fig. 3(a) and (b) as a projection on the (100) and (010) plane respectively. It shows that one molecule is very close to another related to it by the translation *b*. The distance between the carbon atom C(15) of the methoxyl group of one molecule and the oxygen atom O(1) of the other molecule is 2.930 ± 0.003 Å.

Short CH₃...O contacts in crystals are not unusual; a number of them have been reviewed in a paper by Sutor (1963). They mainly occur in the crystal structures of heterocyclic biologically important molecules with, in many cases, a conjugated double bond system. Aflatoxin B₂ clearly satisfies this classification. The shortest CH₃...O contacts described in the review by Sutor are 3.00 Å. In the aflatoxin B₂ crystal, the distance is even 0.07 Å shorter. This short contact points to an interaction between the molecules and, because hydrogen atoms are involved, this interaction may be due to hydrogen bonding. Since the molecule is almost parallel to the (100) plane, two molecules are practically coplanar.

Thus, strings of coplanar molecules occur in the direction of the *b* axis. Within these strings, the molecules are possibly interconnected by hydrogen bonds. The strings are stacked in the direction of the *a* axis with an interplanar distance of *a*/2 = 3.463 Å (*cf.* C(1)...C(8) = 3.468 and C(5)...C(5) = 3.490 Å).

Thermal motion

The atomic vibration ellipsoids have been analysed in terms of the rigid-body tensors of translation (**T**), libration (**L**) and screw motion (**S**) according to Schomaker & Trueblood (1968). Their program, available from the ACA library (No. 1) was adapted for use on an IBM 1800 computer. All the 23 heavy atoms were considered and the results of the analysis are given in Table 8. The r.m.s. deviation between the thermal parameters obtained by the least-squares refinement (further referred to as 'observed') and those calculated from the rigid-body tensors is 0.0043 Å². Although this value is not very low, the resulting modes of motion are physically reasonable. The translation tensor **T** is slightly anisotropic, the direction of the smallest vibration being almost parallel to the *a* axis (the angle between the two directions is 4°48'). A possible explanation is that the packing between the molecules is closest in the *a* direction, which is approximately perpendicular to the plane of the molecule. The unique origin obtained by symmetrizing **S** was used to calculate the inertia tensor **I** (principal axes and the cosines of their angles with respect to the crystallographic axes) of the aflatoxin B₂ molecule. There is a striking relation between **L** and **I**, which are both anisotropic. The largest principal axis

of **L** has almost the same direction as the smallest principal axis of **I**. In other words, the libration is largest around an axis, the moment of inertia of which is smallest. The reverse is also valid (see Table 8): the libration is smallest around an axis of which the moment of inertia is largest. The angles between the corresponding axes of both tensors are 2°38', 6°18' and 5°4', respectively. The libration tensor related to an axial system defined by the principal axes of the inertia tensor has also been calculated (Table 8) and is practically diagonal. The distance between the unique origin and the centre of gravity is small: 0.137 Å, compared for example with 0.542 Å for 7,7,8,8-tetracyanoquinodimethane (Hanson, 1968*a*). There are relatively large differences between observed and calculated *U*₁₁'s (r.m.s. Δ*U*₁₁ = 0.0091 Å²), as can be seen in Table 9. Especially, the values found for the atoms O(1) and O(5) are too large compared with the calculated values, indicating a possibly non-rigid thermal motion of these atoms perpendicular to the plane of the molecule.

Table 8. Rigid-body thermal parameters

$\mathbf{T} = \begin{pmatrix} 193 & -4 & 8 \\ & 247 & 0 \\ & & 273 \end{pmatrix} \times 10^{-4} \text{ \AA}^2$			
$\sigma(\mathbf{T}) = \begin{pmatrix} 17 & 13 & 11 \\ & 12 & 10 \\ & & 10 \end{pmatrix} \times 10^{-4} \text{ \AA}^2$			
$\mathbf{L} = \begin{pmatrix} 37 & -5 & 5 \\ & 93 & -24 \\ & & 142 \end{pmatrix} \times 10^{-1} (^{\circ})^2$			
$\sigma(\mathbf{L}) = \begin{pmatrix} 4 & 5 & 5 \\ & 6 & 5 \\ & & 8 \end{pmatrix} \times 10^{-1} (^{\circ})^2$			
Unique origin (Å)*	-0.336	4.885	15.941
Centre of gravity (Å)	-0.308	4.976	15.866
Principal axes of I †	Direction cosines (× 10 ⁴)		
Eigenvalue			
1219 Å ²	831	-3283	9409
3268	-418	-9445	-3259
4318	9957	-122	-922
Principal axes of T ‡	Direction cosines (× 10 ⁴)		
Eigenvalue			
0.0271 Å ²	-659	1707	-9831
0.0221	-517	9833	1742
0.0155	9965	616	-564
Principal axes of L	Direction cosines (× 10 ⁴)		
Eigenvalue			
15.20 (°) ²	594	-3767	9243
8.34	514	-9236	-3798
3.70	9969	694	-359
Libration tensor L related to the principal axis of I			
$\mathbf{L} = \begin{pmatrix} 46 & 1 & 0 \\ & 25 & 1 \\ & & 11 \end{pmatrix} \times 10^{-4} \text{ rad}^2$			

r.m.s. difference between observed and calculated *U*_{ij}: 43 × 10⁻⁴ Å².

* This origin symmetrizes **S**.

† Calculated for the unique origin, using atomic weights instead of mass weights.

‡ Calculated after **S** had been symmetrized.

Table 9. Differences between observed and calculated U_{ij} ($\times 10^4 \text{ \AA}^2$)

	U_{11}	U_{22}	U_{33}	U_{12}	U_{13}	U_{23}
C(1)	-117	+4	+7	+10	+6	-2
C(2)	-128	+7	-23	-2	+20	+1
C(3)	-23	-7	-26	-15	-8	+4
C(4)	-8	+24	-9	+1	+1	-4
C(5)	+37	+0	+0	+19	+9	-12
C(6)	-107	+25	-11	-9	-46	+13
C(7)	+108	+6	+0	+2	-13	+12
C(8)	+1	+16	+13	+8	-28	+3
C(9)	+11	-3	+8	+15	-17	+15
C(10)	-50	-4	+6	-21	-26	-2
C(11)	-24	+30	+12	+3	+8	-17
C(12)	-80	+28	-10	-27	-19	+32
C(13)	-36	-10	-4	+58	+55	+7
C(14)	+3	-17	-12	-0	+55	-9
C(15)	-105	-22	+30	+2	+31	-3
C(16)	+85	-34	+9	-32	+26	-4
C(17)	+22	+105	+70	+32	+15	-66
O(1)	+207	-3	+5	+4	-0	+1
O(2)	-68	-20	-0	+2	-29	+33
O(3)	+52	-44	-28	+21	+3	+2
O(4)	+8	-12	-32	-51	-1	-0
O(5)	+237	-30	-31	+5	-18	-1
O(6)	-24	-37	+26	-29	-15	-4
r.m.s. ΔU_{ij}	91	30	20	22	25	18

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